

THERMAL EXPANSION OF HEAVY-FERMION SUPERCONDUCTORS

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We have measured the coefficients of linear thermal expansion ( $\alpha$ ) of single-crystalline samples of the heavy-fermion superconductors  $UPt_3$ ,  $URu_2Si_2$  and  $UBe_{13}$ . For the non-cubic systems  $\alpha$  shows a remarkable anisotropy. On entering the superconducting state a strong reduction of the volume expansion is observed for all three systems.

1. INTRODUCTION

The occurrence of superconductivity in the heavy-fermion compounds is undoubtedly an intriguing phenomenon and has received a wide interest, because of the unusual behaviour in the superconducting state. The measured large anomalies in the specific heat at the superconducting transition temperature ( $T_c$ ) have shown unambiguously that the heavy fermions take part in the superconducting condensate. It is then of interest to investigate the corresponding volume and shape effects. Our normal state volume thermal expansion ( $\alpha_v$ ) measurements (1-3) yield large linear electronic terms at low temperatures,  $\alpha_v = aT$ , just as the specific heat,  $c = \gamma T$ , and give rise to large electronic Grüneisen parameters,  $\Gamma \propto a/\gamma \approx 50$  (4,5). From this we conclude that large volume changes can be expected at  $T_c$ . In this paper we present thermal-expansion measurements on the uranium based heavy-fermion compounds  $UPt_3$ ,  $URu_2Si_2$  and  $UBe_{13}$ . In order to study the anisotropy in the coefficients of linear thermal expansion,  $\alpha = L^{-1}(dL/dT)$ , the experiments have been performed on single-crystalline specimens.

2. EXPERIMENTAL

The samples were mounted in a dilatation cell (6) machined of OFHC copper. The thermal expansion was measured using a sensitive 3-terminal capacitance method with a detection limit of 0.1 Å. In the case of  $UPt_3$ , the cell was attached to the mixing chamber of a dilution refrigerator, whereas in the case of  $URu_2Si_2$  and  $UBe_{13}$  it was attached to the cold plate of a  $^3He$  cryostat. Data points were gathered stepwise ( $\Delta T > 5$  mK) in order to ensure thermal equilibrium of sample and cell. The data have been corrected for the cell effect, i.e. the signal of the cell with a copper sample.

3. RESULTS

The thermal expansion of hexagonal  $UPt_3$  has been studied for two samples (7). In Fig.1 we present data taken on the specimen prepared by L.Taillefer (size  $1 \times 2 \times 3$  mm<sup>3</sup>). The superconducting transition is most clearly observed in the expansion along the hexagonal axis ( $\alpha_H$ ), where a negative jump occurs at a (midpoint) transition temperature of 490 mK.

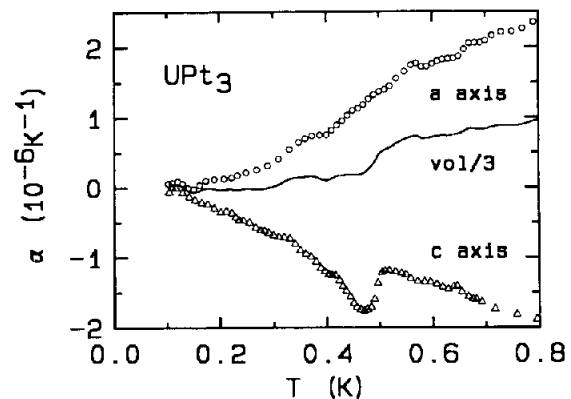


FIGURE 1  
 Coefficient of linear thermal expansion of  $UPt_3$  for the a and c axis as indicated. The solid line represents  $\alpha_v/3$ .

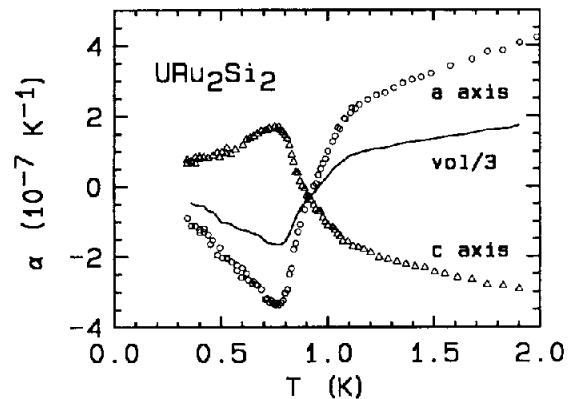


FIGURE 2  
 Coefficient of linear thermal expansion of  $URu_2Si_2$  for the a and c axis as indicated. The solid line represents  $\alpha_v/3$ .

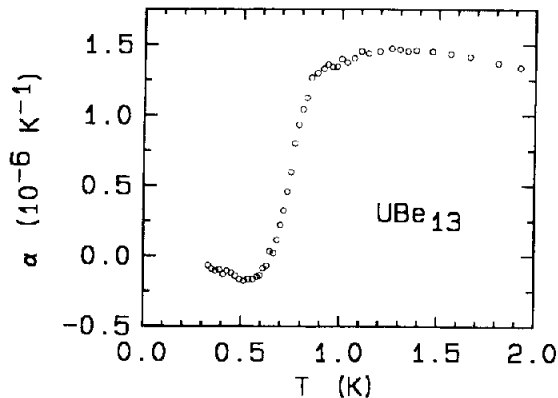


FIGURE 3  
Coefficient of linear thermal expansion of  $\text{UBe}_{13}$ .

For the coefficient of thermal expansion in the basal plane ( $\alpha_{\perp}$ ) no discontinuity is observed at  $T_c$ . The temperature variation of  $\alpha_{\perp}$  changes from linear above  $T_c$  to quadratic there below. The overall anisotropy is preserved below  $T_c$ : the c-axis contracts and the basal plane expands with raising temperature. Recently, specific heat measurements have revealed the existence of two transitions at  $T_{c1} = 490$  mK and  $T_{c2} = 430$  mK in  $\text{UPt}_3$  (8), that are thought to originate from a splitting of the superconducting transition by a symmetry breaking field. The transition at  $T_{c2}$  could not be detected from the data in Fig.1, but having increased the temperature resolution and the sensitivity of our dilatometer we have indeed resolved an anomaly at  $T_{c2}$  (9).

The coefficients of linear thermal expansion along ( $\alpha_{\parallel}$ ) and at right angles ( $\alpha_{\perp}$ ) to the tetragonal axis of  $\text{URu}_2\text{Si}_2$  have been measured in the temperature range 0.3-2 K (see Fig.2). The sample (size  $5 \times 5 \times 5$  mm<sup>3</sup>) has been prepared by the Czochralski technique by A.A.Menovsky. Superconductivity sets in near 1.1 K, below which temperature broad anomalies are observed:  $\alpha_{\perp}$  shows a rapid drop below  $T_c$  and attains a negative minimum at 0.75 K, while  $\alpha_{\parallel}$  increases and has a positive maximum at 0.75 K. We would like to draw special attention to the reversal of anisotropy: well below  $T_c$  the tetragonal axis expands and the basal plane contracts, whereas in the normal state the contrary takes place.

The coefficient of linear thermal expansion of cubic  $\text{UBe}_{13}$  (prepared by J.L.Smith; sample size  $2 \times 4 \times 6$  mm<sup>3</sup>) is shown in Fig.3, for  $0.3 < T < 2$  K. The normal-state expansion shows a weak maximum near 1.3 K, that is attributed to the Kondo-lattice effect. Superconductivity sets in at 0.85 K, where  $\alpha$  starts to fall rapidly. A negative minimum is attained at 0.52 K. Similar data have been obtained by Ott (10).

For  $\text{UPt}_3$  and  $\text{URu}_2\text{Si}_2$  the volume expansion has been calculated from  $\alpha_v = \alpha_{\parallel} + 2\alpha_{\perp}$  (see Figs 1 and 2). For all three compounds we observe a strong reduction of  $\alpha_v$  on entering the superconducting state, which indicates that a positive volume is associated with the formation of the ground state. The Grüneisen parameters drop from large positive values above  $T_c$  to small negative values well below  $T_c$ .

#### 4. DISCUSSION

In a first analysis the anomaly in  $\alpha_v$  at  $T_c$  can be used to calculate the pressure dependence of  $T_c$  via the Ehrenfest relation:  $dT_c/dP = V_m T_c \Delta\alpha_v / \Delta c_p$  ( $V_m$  is the molar volume, and  $\Delta\alpha_v$  and  $\Delta c_p$  are the jumps in the volume expansion and specific heat at  $T_c$ , respectively). From the data in Figs 1-3 and published specific heat data we derive values for  $dT_c/dP$  of -16.3, -101 and -19.4 mK/kbar for  $\text{UPt}_3$ ,  $\text{URu}_2\text{Si}_2$  and  $\text{UBe}_{13}$ , respectively. The values for the depression of  $T_c$  can be compared with the ones obtained from resistivity measurements as function of pressure: -24 (11), -95 (12) and -16 mK/kbar (13), respectively.

The anisotropy in  $\alpha$  near  $T_c$  can be used to determine the uniaxial stress effects on  $T_c$ . For instance, in the case of  $\text{UPt}_3$  the anomaly at  $T_c$  for  $\alpha_{\parallel}$  indicates a decrease of  $T_c$  for uniaxial stress along the c-axis, whereas stress in the basal plane will hardly effect  $T_c$  (7). In the case of  $\text{URu}_2\text{Si}_2$  stress along c will increase  $T_c$ , whereas stress applied in the basal plane will lead to a reduction of  $T_c$ . In order to study further the volume and shape effects at the superconducting transitions in the heavy-fermion superconductors thermal expansion measurements in an external field are in progress.

#### ACKNOWLEDGEMENTS

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