THE UPPER CRITICAL FIELD OF SUPERCONDUCTING UPt3 ALLOYED WITH BORON.

T. VORENKAMP, A. DE VISSER, R. WESTER, A.A. MENOVSKY and J.J.M. FRANSE.

Natuurkundig Laboratorium der Universiteit van Amsterdam, Valckenierstraat 65, 1018 XE, Amsterdam, The Netherlands.

Measurements of  $H_{c2}(T)$ , on two very pure samples of hexagonal UPt3B<sub>0.11</sub>, are presented for H//a, H//b and H//c. A "kink" in  $H_{c2}(T)$  is observed for both basal plane directions. The values of  $dH_{c2}/dT$  are comparable with those of pure UPt3.

# 1. INTRODUCTION

Heavy fermion superconductors attract a lot of interest because the superconducting order parameter in these compounds might be nonscalar, leading to the possible existence of more than one superconducting phase. In particular for UPt3 there is growing evidence for a multi-component phase diagram. Specific heat measurements in zero field of UPt3 (1) reveal two anomalies separated by 60 mK, suggesting that there are two nearly degenerate superconducting states. The upper critical field of UPt3 shows large anisotropy when the field direction is rotated from the basal plane to the hexagonal axis (c-axis), and a change of slope is seen for a field direction in the basal plane (2-6) This change in slope is reported to be a sharp kink by some authors (3-5), and a positive curvature by others (2,6). Recently a theoretical approach is made (7-9), explaining the double peak in the specific heat with a coupling between the superconducting order parameter and the observed small antiferromagnetic moment Ms. In this model  $M_{\rm S}$  is fixed to one basal plane direction and acts as a symmetry breaking field. Calculations predict a kink in  $H_{c2}(T)$ for one direction in the basal plane, and anisotropy for field directions within the basal plane. No basal plane anisotropy has been observed by the authors of ref. (4), whereas in ref. (10) a small anisotropy has been reported. Previously it has been shown that alloying UPta with small amounts of boron leads to an increase of  $T_{\rm c}$  (11). Moreover it has been shown that the double peak in the zero field specific heat is not altered or depressed by the addition of small amounts of boron (12). In this paper we report on  $H_{c2}(T)$  measurements of UPt3B0.11 (nominal concentration of boron) for field directions parallel to the a,b and c-axis. A comparison is made with the data for UPt3. In our investigation special attention has been given to the question whether there is basal plane anisotropy, whether both H//a and H//b show a change in slope, and whether the

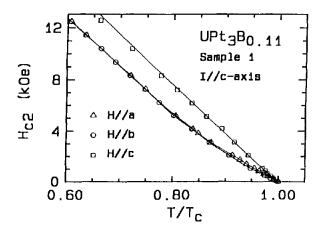
change in slope is a sharp kink or a positive curvature.

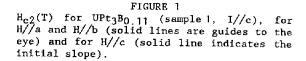
### 2. EXPERIMENTAL

The two samples investigated are both cut same Czochralski-grown single from the crystalline rod. Both samples are annealed in the same way. Growing and annealing procedures described previously (12). homogeneity of the boron content was confirmed by microprobe analysis. The absolute value of the boron content, however, can not be checked in this way. Therefore we will refer to the 11% boron content as a nominal value. The samples are of high purity as is indicated by the low  $\rho_{\rm O}$ -value of 0.29  $\mu\Omega{\rm cm}$  (for I//c-axis). The RRR-value as defined by R(294K)/R(0K) is 536, and the transition temperature in zero field is 559 mK. Sample 1 exhibits a small mosaik of 30, and sample 2 of 50, as was shown by Lauepictures. Resistance was measured in a <sup>3</sup>Hecryostat with a four point AC-method, using a transformer coupled bridge and a current density of 0.2 A/cm<sup>2</sup> (sample 1) and 0.1 A/cm<sup>2</sup> (sample 2). The temperature was measured with a RuO2 thermometer which was calibrated in field against a capacitance thermometer. The samples were cut into the shape of a rectangular bar and the demagnetization correction is estimated to be always less than

## 3. RESULTS AND DISCUSSION

The results are presented in figs 1 and 2 for sample 1 (I//c) and sample 2 (I//b), respectively. The transition temperatures have been determined from the 50% value of the resistance in the normal state. The latter being defined as the point where the resistance deviates from the normal state  $\mathbf{T}^2$  behaviour. The width of the transition as determined from the 10% and 90% values of the normal state resistance is (16±2) mK and does not broaden with magnetic field up to 1.4 Tesla. Refering to the data in figs. 1 and 2, we stress 4 important facts.





- 1) The addition of boron does not appreciably change the characteristics of  $\rm H_{c2}(T)$ . For the basal plane field direction (sample 2)  $\rm dH_{c2}/dT$  is -67 kOe/K below  $\rm T_H/T_c$  = 0.81 and -45 kOe/K there above ( $\rm T_H$  being the temperature where the change in slope occurs). For H//c  $\rm dH_{c2}/dT$  is -69 kOe/K. These values are comparible to the values -64, -44 and -71 kOe/K, respectively, for pure UPt3 (10). 2) For both H//a and H//b  $\rm H_{c2}(T)$  shows a
- 2) For both H//a and H//b  $\rm H_{c2}(T)$  shows a change of slope which sets in at  $\rm T_H/T_c=0.81$  (fig. 1). This implies that the overal behaviour of  $\rm H_{c2}(T)$  in the basal plane is isotropic, and that the earliest theoretical models, using a large in-plane anisotropy of the antiferromagnetic order, need some refinement.
- 3) Nevertheless, from the data in fig. 1 a small in-plane anisotropy remains visible. When we denote this as an increase of  $T_{\rm C}$  by going from the H//b phase line to H//a phase line, the upper limit for the anisotropy is 5 mK for T>T\_H and 2.5 mK for T<T\_H This is in rough agreement with the findings of ref. (10).
- 4) The experimental results for sample 2 (H//b//I) indeed show a sharp kink at  $T_H/T_C=0.81$ , whereas for sample 1  $(H//b \bot I)$  the change of slope sets in above  $T_H/T_C=0.81$  and a small positive curvature is observed. The deviation between the curved phase line of sample 1 and the linear phase line of sample 2 is small, the maximum deviation being 5 mK. The difference between a sharp change of slope or a gradual change of slope (or less sharp kink) might be related to the difference in current densities (10), the current density being a

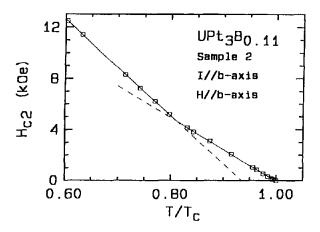


FIGURE 2  $\rm H_{c2}(T)$  for UPt3B0.11 (sample 2, I//b) for H//b (solid lines indicate the two different slopes dH<sub>c2</sub>/dT).

factor two smaller in case of the sharp kink. In conclusion we have presented data for  $\rm H_{\rm c2}(T)$  of single-crystalline UPt3B\_0.11 and commented on the observed anisotropies.

### ACKNOWLEDGEMENTS

The work of T.V. was part of the research program of the foundation FOM. The work of A.d.V. has been made possible by a fellowship of the Royal Netherlands Academy of Arts and Sciences.

# REFERENCES

- (1) R.A. Fisher et al., Phys. Rev. Lett. 62 (1989) 1411.
- (2) J.W. Chen et al., Phys. Rev. B 30 (1984) 1583.
- (3) U. Rauchschwalbe et al., Z. Phys. B Cond. Matt. 60 (1985) 379
- (4) B.S. Shivaram et al., Phys. Rev. Lett. 57 (1986) 1259. B.S. Shivaram et al., Phys. Rev. Lett. 63 (1989) 1723.
- (5) L. Taillefer et al., Physica C 153-155 (1988) 451
- (6) W.K. KwoK, L.E. DeLong, G.W. Crabtree, D.G. Hinks and R. Joynt, (to be published).
- (7) D.W. Hess et al., J. Phys. Condens. Matter 1 (1989) 8135.
- (8) K. Machida et al., J. Phys. Soc. Jpn. 58 (1989) 4116.
- (9) S.K. Sundaram and R. Joynt, Phys. Rev. B 40 (1989) 8780.
- (10) L. Taillefer, Physica B 'in print'.
- (11) T. Vorenkamp et al., J. Magn. Magn. Mat. 76&77 (1988) 531.
- (12) T. Vorenkamp, et al., Physica B 'in print'.